

### Introduction

We have analysed CRM U500 on the **Phoenix TIMS** using Faraday collectors equipped with **Xact** amplifier boards carrying  $1e^{11} \Omega$  resistors and  $1e^{12} \Omega$  (1T  $\Omega$ ) resistors. The 1T  $\Omega$  resistors should in theory offer three times less noise than the  $1e^{11} \Omega$  resistors and this should result in a higher precision for smaller ion beams.

### Noise level

Table 1 and Figure 1 show the noise levels of  $1e^{11} \Omega$  and 1T resistors in the same amplifier housing. Temperature conditions are the same, as the housing is Peltier cooled to 16°C and evacuated using a scroll pump. The  $1e^{11} \Omega$  resistor noise decreases with increasing integration time. This has been demonstrated in Technical note [G10712](#) which also showed that the noise for  $1e^{11} \Omega$  resistors, closely followed the theoretical Johnson Noise. The 1T resistors show lower noise than the  $1e^{11} \Omega$  resistors irrespective of the integration time. However, the noise levels are not three times better than the  $1e^{11} \Omega$  resistors. Typically the noise is better by a factor of 1.5, though the difference is more pronounced for shorter integration times.

This lower than expected improvement in noise has also been reported by Richter et. al 2011 on a different TIMS instrument. For a typical integration time of 10 seconds the noise level on 1T resistors is  $6.5e^{-17}$  A which can be achieved on  $1e^{11} \Omega$  resistors if the noise is integrated for 40 seconds or more. This would mean that at least four 10-second integrations can be made with 1T resistors at the same noise level as one 40-second integration with  $1e^{11} \Omega$  resistor.

### 1T and $1e^{11} \Omega$ Resistor Time Response (Tau)

In the past 1T resistors have demonstrated an extremely slow time response, i.e the time taken for a signal to decay to baseline can be several tens of seconds. Richter et. al 2011, state an idle time of 15 seconds is required before data is collected following a mass jump. This would clearly make measurements involving peak jumping prohibitively slow, especially for small samples.

Figure 2 shows the resistor response for a standard  $1e^{11} \Omega$  resistor, and shows that baseline can be achieved in about one second. Figure 3 shows the tau trace for four 1T resistors. Baseline is achieved between three and four seconds, so they are approximately three times slower than  $1e^{11} \Omega$  resistors. This has also been noted by Tuttas et. al.2007<sup>3</sup>.

Integration time (secs)	Noise (Amps)		
	$1e^{12} \Omega$ n=5	$1e^{11} \Omega$ n=4	$1e^{11}/1e^{12}$
1	2.3E-16	3.8E-16	1.7
5	1.0E-16	1.7E-16	1.6
10	6.5E-17	1.2E-16	1.8
20	5.1E-17	7.3E-17	1.4
40	3.6E-17	6.3E-17	1.7
60	3.6E-17	5.1E-17	1.4
300	3.5E-17	4.9E-17	1.4

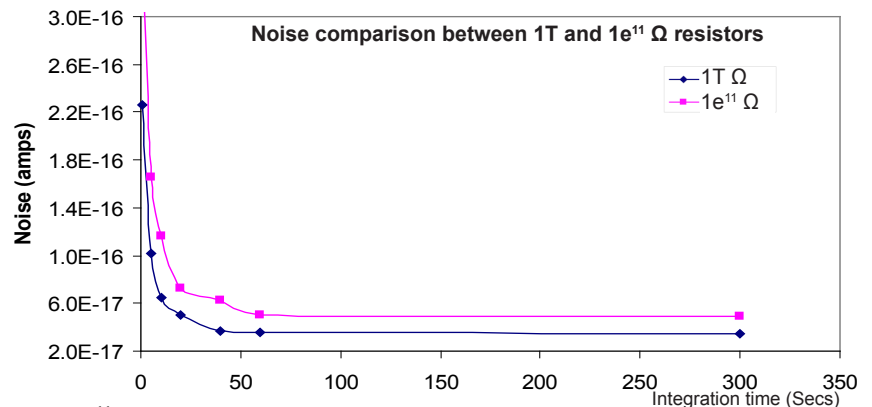


Table 1. & Figure 1. Noise comparison between 1T and  $1e^{11} \Omega$  resistors.

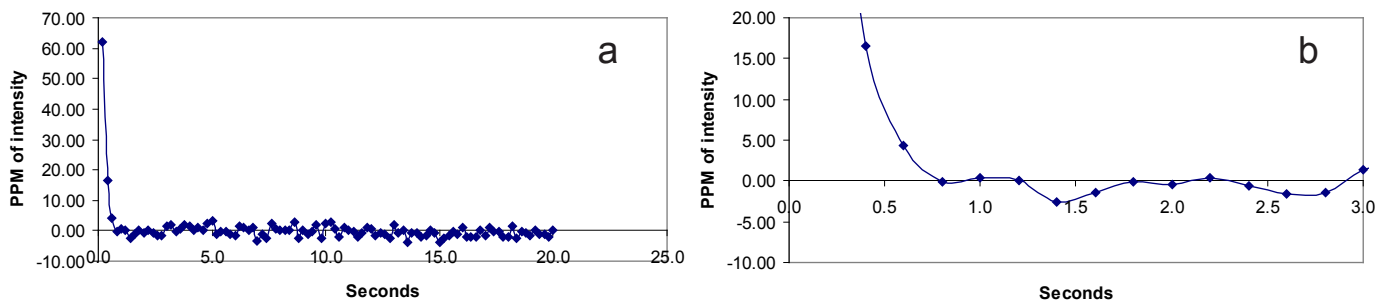


Figure 2. Tau response of a standard Xact  $1e^{11} \Omega$  resistor showing response over 20 seconds (a) and expanded portion covering the first three seconds (b). Baseline is reached in less than 1 second. 5ppm of peak intensity is reached within 0.5 seconds. The Xact amplifier boards are approximately 4 times faster than the previous generation of boards.

### CRM U500

The CRM U500 is particularly useful in demonstrating the performance difference between the two types of resistors.

The major isotopes are equivalent in size ( $^{235}\text{U}/^{238}\text{U} = 0.9997$ ) and can be used to fractionation correct the minor isotope ratios  $^{234}\text{U}/^{238}\text{U}$  and  $^{236}\text{U}/^{238}\text{U}$  whose ratios are 0.010425 and 0.001522 respectively<sup>1</sup>. This large dynamic range in minor/major isotope ratio of two and three orders of magnitude are good indicators of the ability to measure small ion signals.

### Method

2µg loads of U500 were loaded on degassed triple filament assemblies. During analysis the ion signal was maintained at  $6.5 \text{ e}^{-12} \text{ A}$  equivalent to 650mv using a  $1\text{e}^{11} \Omega$  resistor. The relatively low intensity used for these analyses was used to compare the performance of the two types of resistors to analyze small ion beams.

Data was collected in the static mode. A total of 400 ten second integrations were made, in 20 blocks of 20 cycles. Baselines were taken at the start of each block at  $\pm 0.5 \text{ amu}$  for 30 seconds at each mass position. A four second delay was made following each mass jump before integrating the signal or baseline as this is the length of time for the signal to reach baseline with the 1T resistors.

### 1e<sup>11</sup> Ω CRM U500 Faraday data

The Data for  $1\text{e}^{11} \Omega$  are shown in Table 2. The columns in blue are the measured isotope compositions. The mean  $^{234}\text{U}/^{238}\text{U}$  is 0.010419 $\pm$ 0.1%, the mean  $^{235}\text{U}/^{238}\text{U}$  is 0.99994 $\pm$ 0.07% and the mean  $^{236}\text{U}/^{238}\text{U}$  is 0.001516 $\pm$ 0.39%. As expected the precision gets worse with the involvement of the smaller ion

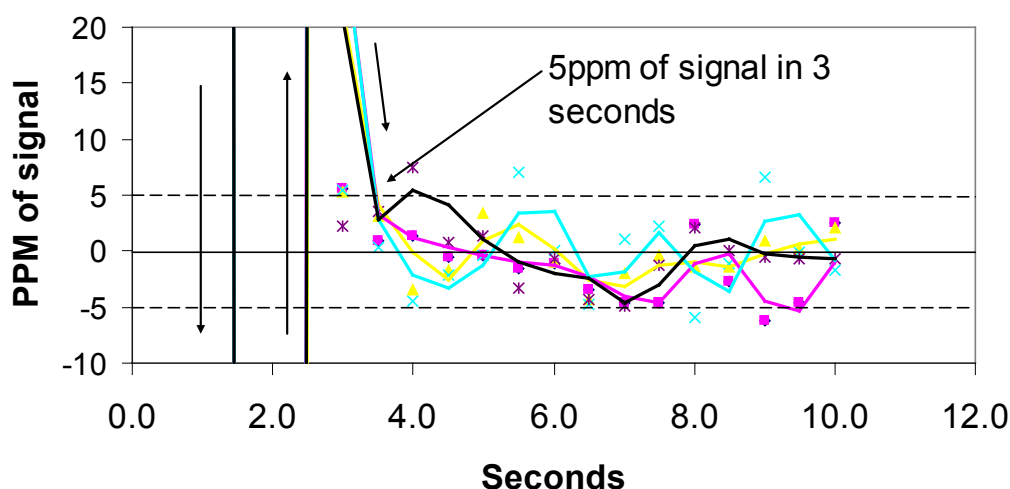
beam, however, there remains an element of non-corrected mass fractionation which means that the reproducibility of the  $^{234}\text{U}/^{238}\text{U}$  is similar to the  $^{235}\text{U}/^{238}\text{U}$ . The mass fractionation effect can be overcome if the  $^{234}\text{U}/^{238}\text{U}$  and  $^{236}\text{U}/^{238}\text{U}$  are fractionation corrected with  $^{235}\text{U}/^{238}\text{U}$ . This is shown in the yellow columns.

The average internal precision (%1se) of  $^{234}\text{U}/^{238}\text{U}$  is 0.01%, and  $^{236}\text{U}/^{238}\text{U}$  0.06%. This is consistent with the  $^{236}\text{U}$  being 6 times smaller than the  $^{234}\text{U}$ . The reproducibility for the  $^{234}\text{U}/^{238}\text{U}$  at 0.06% is also six times better than the 0.36% for the  $^{236}\text{U}/^{238}\text{U}$ . We can use the  $^{234}\text{U}/^{238}\text{U}$  to fractionation correct the  $^{235}\text{U}/^{238}\text{U}$  and a mean ratio of 1.0004 $\pm$ 0.046% 1RSD is obtained, this is shown in the green columns. The reproducibility for the  $^{235}\text{U}/^{238}\text{U}$  is a factor of two larger than the individual measurement precision.

### 1T Ω CRM U500 Faraday data

Data for amplifier boards fitted with 1T resistors are shown in Table 3. For the non-fractionation corrected ratios, the mean  $^{234}\text{U}/^{238}\text{U}$  is 0.010442 $\pm$ 0.16%, the mean  $^{235}\text{U}/^{238}\text{U}$  is 1.00091 $\pm$ 0.12% and the mean  $^{236}\text{U}/^{238}\text{U}$  is 0.001522 $\pm$ 0.14%. Reproducibilities are all the same unlike the  $1\text{e}^{11} \Omega$  data. There is however a startling improvement in the fractionation corrected data shown in the yellow columns.  $^{234}\text{U}/^{238}\text{U}$  is 0.010425 $\pm$ 0.01% and  $^{236}\text{U}/^{238}\text{U}$  is 0.001520 $\pm$ 0.09%. This indicates that the fractionation correction has improved the data by more than an order of magnitude. The normalised  $^{235}\text{U}/^{238}\text{U}$  has a mean ratio of 0.99972 $\pm$ 0.01% 1RSD shown in the green columns. The reproducibility for the  $^{235}\text{U}/^{238}\text{U}$  is a factor of two larger than the individual measurement precisions, and is an order of magnitude better than the uncorrected data.

### 1T signal response



**Figure 3.** Tau response of 1T resistor. At 0 seconds a 5 volt ion beam is switched off the peak by 0.5 amu. The signal decays rapidly and overshoots the baseline after 1 second, it then rebounds at 2.5 seconds, before finally decaying to 5ppm of the 5 volt beam after 3 seconds. This graph shows the response of 4 different collectors.

### Comparison between 1T and 1e<sup>11</sup> Ω.

The isotope data show that there is no difference in isotope composition or precision between the 1T and the 1e<sup>11</sup> Ω resistors for the un-normalized data sets. This is surprising considering that the minor isotope intensities <sup>234</sup>U and <sup>236</sup>U are only 6e<sup>-14</sup> A and 1e<sup>-14</sup> A respectively. The biggest source of error appears to be the fractionation behaviour on the filament. For the normalized ratios the <sup>234</sup>U/<sup>238</sup>U and <sup>236</sup>U/<sup>238</sup>U individual analysis precision is improved by a factor of two with the 1T resistors (Table 4).

Table 4 summarizes the normalized <sup>234</sup>U/<sup>238</sup>U and <sup>236</sup>U/<sup>238</sup>U data for individual analyses. We can calculate the resistor noise by taking the intensity of the minor isotopes and dividing it by the standard deviation of the error. In this case the standard deviation is the standard error multiplied by 20, where 20 is the square root of the number of ratios, which in this case is 400. The noise is approximately 6e<sup>-17</sup> A for the 1T resistor and 1.2e<sup>-16</sup> A for the 1e<sup>11</sup> Ω resistor. This is consistent with the

measured noise for a 10 second integration shown in Table 1.

The reproducibility is better with the 1T resistors by a factor of 6 for the <sup>234</sup>U/<sup>238</sup>U and by a factor of 4 with the <sup>236</sup>U/<sup>238</sup>U. The <sup>235</sup>U/<sup>238</sup>U internal precision is improved by a factor of two while the reproducibility is improved by a factor of 4.5 (Figure 3).

### Conclusion

- 1T resistors are between 1.5 and 2 times quieter than 1e<sup>11</sup> Ω resistors for the same integration time.
- The 1T resistors are approximately 3 times slower than the 1e<sup>11</sup> Ω resistors and they reach baseline in 3-4 seconds.
- There is no difference between the two resistor types for non-fractionation corrected data. Fractionation corrected data shows a factor of 2 increase in precision for ratios involving the minor uranium isotopes.
- Reproducibility is dramatically improved by at least a factor of 5 with the 1T resistors

Filament	<sup>238</sup> U	Measured ratios			Normalised to <sup>235</sup> U/ <sup>238</sup> U = 0.9997				Normalised to <sup>234</sup> U/ <sup>238</sup> U = 0.010425	
		<sup>234</sup> U/ <sup>238</sup> U	<sup>235</sup> U/ <sup>238</sup> U	<sup>236</sup> U/ <sup>238</sup> U	<sup>234</sup> U/ <sup>238</sup> U (N)	%1se	<sup>236</sup> U/ <sup>238</sup> U (N)	%1se	<sup>235</sup> U/ <sup>238</sup> U (N)	%1se
1	6.5E-12	0.010421	1.00061	0.001517	0.010409	0.008	0.001516	0.04	1.00087	0.006
2	6.7E-12	0.010418	0.99933	0.001527	0.010423	0.010	0.001525	0.07	0.99983	0.008
3	6.1E-12	0.010400	0.99903	0.001512	0.010408	0.011	0.001512	0.08	1.00084	0.009
4	6.5E-12	0.010419	0.99960	0.001518	0.010420	0.012	0.001518	0.05	1.00006	0.009
5	6.5E-12	0.010431	1.00034	0.001516	0.010421	0.011	0.001515	0.05	0.99994	0.008
6	6.6E-12	0.010429	1.00045	0.001518	0.010417	0.011	0.001517	0.05	1.00016	0.009
7	6.8E-12	0.010432	1.00058	0.001519	0.010419	0.009	0.001518	0.04	1.00008	0.007
8	6.1E-12	0.010405	0.99944	0.001506	0.010408	0.014	0.001507	0.06	1.00090	0.010
9	6.5E-12	0.010421	1.00076	0.001508	0.010405	0.012	0.001507	0.05	1.00106	0.009
10	6.3E-12	0.010411	0.99930	0.001518	0.010416	0.011	0.001518	0.05	1.00030	0.008
<b>MEAN</b>	6.5E-12	<b>0.010419</b>	<b>0.99994</b>	<b>0.001516</b>	<b>0.010415</b>	<b>0.01</b>	<b>0.001515</b>	<b>0.06</b>	<b>1.00040</b>	<b>0.008</b>
<b>1SD</b>	2.2E-13	<b>0.000011</b>	<b>0.000066</b>	<b>0.000006</b>	<b>0.000006</b>		<b>0.000006</b>		<b>0.00046</b>	
<b>%1 RSD</b>	3.4	<b>0.10</b>	<b>0.07</b>	<b>0.39</b>	<b>0.06</b>		<b>0.36</b>		<b>0.046</b>	

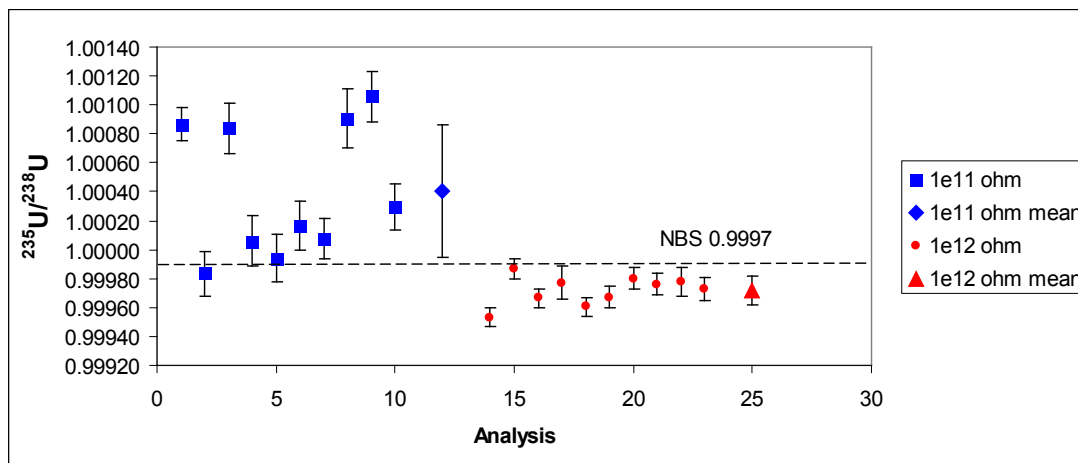
Table 2. Isotope ratio performance using 1e<sup>11</sup> Ω resistors.

Filament	<sup>238</sup> U	Measured ratios			Normalised to <sup>235</sup> U/ <sup>238</sup> U = 0.9997				Normalised to <sup>234</sup> U/ <sup>238</sup> U = 0.010425	
		<sup>234</sup> U/ <sup>238</sup> U	<sup>235</sup> U/ <sup>238</sup> U	<sup>236</sup> U/ <sup>238</sup> U	<sup>234</sup> U/ <sup>238</sup> U (N)	%1se	<sup>236</sup> U/ <sup>238</sup> U (N)	%1se	<sup>235</sup> U/ <sup>238</sup> U (N)	%1se
1	6.0E-12	0.010437	1.00041	0.001524	0.010427	0.005	0.001523	0.02	0.999532	0.003
2	5.9E-12	0.010421	0.99956	0.001520	0.010422	0.005	0.001520	0.02	0.999872	0.004
3	6.0E-12	0.010448	1.00133	0.001522	0.010425	0.005	0.001520	0.02	0.999666	0.003
4	6.0E-12	0.010436	1.00055	0.001520	0.010424	0.008	0.001519	0.05	0.999770	0.006
5	6.0E-12	0.010467	1.00265	0.001524	0.010427	0.005	0.001521	0.02	0.999606	0.003
6	6.1E-12	0.010449	1.00140	0.001523	0.010425	0.005	0.001521	0.02	0.999670	0.004
7	6.0E-12	0.010455	1.00194	0.001523	0.010423	0.005	0.001520	0.03	0.999800	0.004
8	6.1E-12	0.010460	1.00228	0.001523	0.010426	0.005	0.001521	0.02	0.999761	0.004
9	5.7E-12	0.010421	0.99952	0.001519	0.010425	0.007	0.001519	0.03	0.999779	0.005
10	5.7E-12	0.010421	0.99947	0.001519	0.010424	0.005	0.001519	0.03	0.999726	0.004
<b>MEAN</b>	6.0E-12	<b>0.010442</b>	<b>1.00091</b>	<b>0.001522</b>	<b>0.010425</b>	<b>0.01</b>	<b>0.001520</b>	<b>0.03</b>	<b>0.99972</b>	<b>0.004</b>
<b>1SD</b>	1.4E-13	<b>0.000017</b>	<b>0.00118</b>	<b>0.000002</b>	<b>0.000002</b>		<b>0.000001</b>		<b>0.00010</b>	
<b>%1 RSD</b>	2.3	<b>0.16</b>	<b>0.12</b>	<b>0.14</b>	<b>0.01</b>		<b>0.09</b>		<b>0.010</b>	

Table 3. Isotope ratio performance using 1T ohm resistors.

	$^{234}\text{U}/^{238}\text{U}$ (N)	%1se	%1SD	Calculated Noise	$^{236}\text{U}/^{238}\text{U}$ (N)	%1se	%1SD	Calculated Noise
1T	0.010425	0.0053	0.1054	6.6E-17	0.001520	0.027	0.542	4.9E-17
1E+11	0.010415	0.0109	0.2187	1.4E-16	0.001515	0.055	1.109	1.0E-16

**Table 4.** Normalized minor isotope compositions and the calculated noise levels, assuming the noise is entirely from the minor isotope.



**Figure 3.** Comparison between  $^{235}\text{U}/^{238}\text{U}_N$  for CRM U500 using  $1\text{ e}^{11}\ \Omega$  resistor and 1T resistor.  $^{238}\text{U}$  was  $6.5\text{ e}^{-12}\ \text{A}$  for these measurements.

### References

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2. Palacz Z, Wanless D. 2012. Mass Fractionation behaviour of U isotopes during Thermal Ionization, the potential for mass fractionation correction using double spike. Isotopx Application note [G10712](#).
3. Dietmar Tuttas, Johannes Schwieters, Norbert Quaas, Claudia Bouman, 2007. Thermo Fisher Scientific, Bremen, Germany Improvements in TIMS High Precision Isotope Ratio Measurements for sample sizes. Triton Application Note 30136.

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